

# Thermo-elastic properties of dense YSZ and porous Ni-ZrO<sub>2</sub> monolithic and isotropic materials

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Attempts are made to correlate the structure and properties of dense and porous Yttria Stabilized Zirconia to establish optimal thermo-elastic properties for better performance at elevated temperatures. Temperature and compositional dependence of isotropic elastic bulk properties (Young's modulus, Poisson's ratio and Shear and Bulk moduli) are determined using the stiffness constants reported in the literature. Anisotropy of Yttria Stabilized Zirconia increases with composition of Yttria dosage in Zirconia. Optimal composition of 12 mole% Yttria stabilized Zirconia is slightly better than 8 mole% Yttria Stabilized Zirconia based on the improved thermo-elastic properties for better performance at higher temperature (RT-400°C). However 15.5 mole% YSZ seems to be more suitable from the view point of thermo-elastic performance at elevated temperatures (500–800°C). Polycrystalline properties predicted are within 5% error limits of the experimental values for Young's modulus of 12 mole% Yttria Stabilized Zirconia. Numerical prediction of Young's modulus of 12 mole% YSZ for (100) orientation is 362 GPa as compared to experimental value of 370 GPa reported in literature.

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## 1. Introduction

Yttria Stabilized Zirconia (YSZ) is widely used as material for fuel-cell fabrication either by tape casting or sol-gel techniques as an alternative source for power generation. Thin membranes made of dense YSZ often constitute electrolyte material whereas porous NiO/YSZ forms the anode material for the planar design of solid oxide fuel cell. A various types of design strategies are utilized by various research groups [1–7] to develop the optimum alternative source of power by developing fuel-cell batteries for use in automobile or telecommunication industries. The success of design and developments of such solid oxide fuel-cell power by and large depends on well-defined thermal, elastic and plastic properties of the materials varying with temperature. A major stumbling block in the fuel-cell designs using Yttria Stabilized Zirconia materials is the lack of temperature dependant thermo-physical, elastic and plastic properties of dense (YSZ) and porous (Ni/YSZ) materials. A purpose of this paper is to com-

pute the elastic properties viz, elastic moduli, Young's E, Shear -G, Bulk - K and Poisson's ratio,  $\nu$  of dense YSZ and porous Ni-YSZ materials as a function of temperature and composition of Yttria doped to stabilize the cubic crystal structure of the ZrO<sub>2</sub>. The elastic properties of isotropic single crystal or cubic polycrystalline structure of YSZ are calculated here using the elastic stiffness constants by computing the compliance matrix coefficients.

The single crystal elastic constants of cubic YSZ were determined at room temperature from various techniques viz. ultrasonic wave measurements for 8 and 12 mole% Yttria by Pace *et al.* [8], classic vibration by Raleigh's method for 8.8 mole% Yttria by Buckley *et al.* [9] and Brillouin scattering of light between 8 and 20 mole% Yttria by Aleksandrov *et al.* (references: [3–5] in) [10]. The single crystal elastic constants by ultrasonic method for 8 mole% Yttria Stabilized Zirconia at room temperature and also for YSZ phase containing 11.1, 12.1, 15.5 and 17.9 mole% Yttria were measured from 20 to 700°C by

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Kandil *et al.* [10]. The only high temperature elastic modulus of YSZ reported in the literature is  $155 \pm 5$  GPa at  $T = 900^\circ\text{C}$  by Seluc *et al.* [4]. The temperature dependant elastic properties are not available for dense YSZ and especially for porous Ni-YSZ materials. Hence attempts are made to compute the temperature dependant elastic as well the thermal properties of dense and porous materials of Yttria Stabilized Zirconia before and after reduction of NiO-YSZ materials.

## 2. Experimental

The monolithic and multilayer films on Ni-YSZ substrates were commercially prepared by tape-cast technique. The specimens were used to test bending strength of monolithic or multilayers. The specimens were tested in 3-point bending fixture suitably developed for characterization of ceramic materials as per ASTM procedure. The microstructural features and qualitative and quantitative study of various phases, porosity and fractography were performed using JEOL-3 type of SEM. The micro-mechanics involved in flexural failure and its fractographic analysis will be discussed in a separate communication [11].

## 3. Numerical procedure

The stress-strain relations of linearly elastic isotropic materials are given by Hook's law as per Equation 1:

$$\sigma_i = C_{ij} \cdot \varepsilon_j \quad (1a)$$

and

$$\varepsilon_i = S_{ij} \cdot \sigma_j \quad (1b)$$

Where, elastic constants  $C_{ij}$  and  $S_{ij}$  are stiffness and compliance constants of tensorial stress and strain vectors respectively.

A reduced form of Voigt matrix notation of stress ( $\sigma_{ij}$ ) and strain ( $\varepsilon_{ij}$ ) are represented by Equation 2:

$$\begin{Bmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{33} \\ \sigma_{23} \\ \sigma_{31} \\ \sigma_{12} \end{Bmatrix} = [C_{ij}] \cdot \begin{Bmatrix} \varepsilon_{11} \\ \varepsilon_{22} \\ \varepsilon_{33} \\ \varepsilon_{23} \\ \varepsilon_{31} \\ \varepsilon_{12} \end{Bmatrix} \quad (2a)$$

and

$$\begin{Bmatrix} \varepsilon_{11} \\ \varepsilon_{22} \\ \varepsilon_{33} \\ \varepsilon_{23} \\ \varepsilon_{31} \\ \varepsilon_{12} \end{Bmatrix} = [S_{ij}] \cdot \begin{Bmatrix} \sigma_{11} \\ \sigma_{22} \\ \sigma_{33} \\ \sigma_{23} \\ \sigma_{31} \\ \sigma_{12} \end{Bmatrix} \quad (2b)$$

Symmetry of cubic crystal structure of ceramics (e.g, 8YSZ, MgO, NaCl) [8, 12] reduces the  $6 \times 6$  matrix of stiffness (and compliances) coefficients to three independent elastic constants  $C_{11}$ ,  $C_{12}$  and  $C_{44}$  and the remaining 33 coefficients are zero [8, 13, 14]. The corresponding three nonzero compliance constants  $S_{11}$ ,  $S_{12}$  and  $S_{44}$  in cubic crystal are given by Equations 3a–3c:

$$S_{11} = \frac{(C_{11} + C_{12})}{(C_{11} - C_{12})(C_{11} + 2C_{12})} \quad (3a)$$

$$S_{12} = \frac{-C_{12}}{(C_{11} - C_{12})(C_{11} + 2C_{12})} \quad (3b)$$

and

$$S_{44} = \frac{1}{C_{44}} \quad (3c)$$

Isotropic elastic properties of cubic crystals can be defined in terms of the compliance constants. Isotropic Young's modulus can be defined as an inverse of the compliance coefficients,  $S_{11}$  and isotropic shear modulus as half of the inverse of the difference between the compliance constants  $S_{11}$  and  $S_{12}$ . Isotropic bulk modulus can be defined as one third of the inverse of compliance factor ( $S_{11} + 2S_{12}$ ). Three isotropic elastic moduli of the cubic crystals are given by Equation 4a–4c. Poisson's ratio can be defined as the negative ratio of  $S_{12}$  to  $S_{11}$  as represented by Equation 5. Anisotropy of materials can be defined by a factor ( $AF$ ) in Equation 6, if it is close to unity the material is said to be highly isotropic. The properties of isotropic cubic materials are independent of the direction of loading. The cubic material is anisotropic if  $AF > 1$ , elastic properties of the cubic material will then be dependent on the orientation of crystal structure. Young's modulus of a highly anisotropic cubic material depends on direction of external load applied on the material and orientation of the cubic crystal. The directional dependence of elastic modulus of anisotropic cubic material is defined by Cook and Pharr [12] in terms of directional cosine:  $R_{\langle hkl \rangle} = [l_1^2 \cdot l_2^2 + l_2^2 \cdot l_3^2 + l_3^2 \cdot l_1^2]$  by Equation 7. Where,  $l_i$  are direction cosines of axis of loading with respect to  $\langle 100 \rangle$  direction for cubic material.

$$E_{\text{Iso}} = \frac{1}{S_{11}} \quad (4a)$$

$$G_{\text{Iso}} = \frac{1}{2(S_{11} - S_{12})} \quad (4b)$$

and

$$K_{\text{Iso}} = \frac{1}{3(S_{11} + 2S_{12})} \quad (4c)$$

$$\nu = -\left(\frac{S_{12}}{S_{11}}\right) \quad (5)$$

$$AF = \frac{2C_{44}}{(C_{11} - C_{12})} \quad (6)$$

$$E_{(hkl)} = \frac{1}{(S_{11} - [2(S_{11} - S_{12}) - S_{44}] * R_{hkl})} \quad (7)$$

## 4. Results and discussion

The dense Yttria Stabilized Zirconia (YSZ) tape-cast in the form of thin membranes of 25–50  $\mu\text{m}$  are widely used in the fabrication of electrolytes for fuel-cells consisting of multiple electrode assemblies (MEAs). The porous materials such as YSZ reduced to Ni-ZrO<sub>2</sub> for anode and La<sub>(1-x)</sub>Sr<sub>x</sub>MnO<sub>3</sub> or La<sub>(1-x)</sub>Sr<sub>x</sub>CoO<sub>3</sub> in definite proportion (40–50%) of YSZ for cathode are commonly used in the fabrication of MEA stacks of the fuel cell. The porosity of the electrodes in composite ceramic MEA stacks helps to transport the ions to increase the electrical conductivity and reduces the thermal diffusivity due to porosity. The mechanical strength of the composite MEA reduces considerably in comparison with individual strength of dense electrolyte or electrodes or Voigt bound modulus based average bending strength using rule of mixture of composite cell materials. Hence, optimum performance of fuel-cell is a compromise between electrical and thermo-mechanical properties of such materials.

### 4.1. Structure and property correlations

In the following two sections attempt is made to correlate the microstructure features with the thermo-mechanical properties of dense and porous Yttria Stabilized Zirconia (YSZ).

#### 4.1.1. Dense YSZ

The crystal structure of Zirconia doped with 8 mole% or more of Yttria (Y<sub>2</sub>O<sub>3</sub>) predominantly consists of cubic crystalline structure, which exhibits optimum thermal and electrical conductivity and reasonably good mechanical properties. Typical microstructures obtained by using SEM pictures of dense YSZ and that of sintered and reduced NiO-YSZ along with its back scattered image, are shown in Fig. 1. The micro-structural qualitative features and phase quantification were carried out by scanning the back scattered secondary electron beams. SEM microstructure of fully stabilized dense 8YSZ confirmed cubic structure and reduced Ni-YSZ had porous structure. Occasionally pinhole type of porosity, were found scattered even in the dense material with an average pore size of 0.1–1  $\mu\text{m}$ . The pinhole porosity of the dense material could be due to 3% volume expansion of ZrO<sub>2</sub> during sintering (1200–1400°C) and there by resulting in shrinkage porosity of dense material of 3–5% maximum after subsequent cooling to ambient. Obviously 8 mole% of Y<sub>2</sub>O<sub>3</sub> stabilizer dosage in ZrO<sub>2</sub> may be inadequate to avoid the

porosity in YSZ even though it could fully stabilize cubic structure. On the contrary, a limited amount of shrinkage porosity in the nearly dense YSZ material might be beneficial to the ionic transportation through the electrolyte. Effect of composition of Yttria doped in ZrO<sub>2</sub> on ionic conductivity in Fig. 2, clearly indicates that a minimum of 8–9 mole% of Yttria Stabilized Zirconia (YSZ) is necessary for fabrication of thin but dense membranes for fuel-cell electrolytes. The ionic conductivity of electrolyte decreases slightly from 15 to 11 mho/meter (ohm-m)<sup>-1</sup> at 1273°K by increasing Yttria composition from 8 to 12 mole% in YSZ from work of Badwal [15]. The ionic conductivity of 12 mole% YSZ is comparable to that of 8 mole% YSZ at intermediate temperatures (600–770°K). At higher temperatures (900–1111°K) ionic conductivity of 12 mole% YSZ is comparable to that of 3 mole% YSZ [16]. However, interfacial ionic conductivity of typical A<sub>2</sub>B<sub>2</sub>O<sub>7</sub> type of La<sub>2</sub>Zr<sub>2</sub>O<sub>7</sub> compounds formed in composite cathodes (LSM/YSZ) Yttria stabilized Zirconia electrolytes. Interfacial ionic conductivity calculated from the Arrhenius relation (Fig. 2B) is ~1.2E-3 S/m at 1273°K.

The lattice parameter,  $a_0$  of Yttria Stabilized Zirconia over a wide range of dosage of Y<sub>2</sub>O<sub>3</sub> stabilizer depends on density,  $\rho$  of a dense material as per the model of Aleksandrov *et al.* [17] in Equation 8:

$$a_0 = \left[ \left( C - \frac{b \cdot M}{(100 + M)} \right) \cdot (1/\rho) \right]^{\frac{1}{3}} \quad (8)$$

(Where,  $M$ –mole%Y<sub>2</sub>O<sub>3</sub>;  $C$  and  $b$  are material constants for binary system Y<sub>2</sub>O<sub>3</sub>.ZrO<sub>2</sub> given in reference [17])

At constant molar composition of Yttria, density of Yttria Stabilized Zirconia is inversely proportional to cubic exponent of lattice parameter and hence the lattice constant as per Equation (8). Higher densification of YSZ can be achieved by decreasing lattice parameter with increasing Yttria dosage in Zirconia. This is ultimately governed by maximum packing density of the atoms in cubic structure of YSZ.

At higher temperature Yttria Stabilized Zirconia undergoes thermal expansion inducing thermal strains due to a change in lattice parameter or lattice spacing as per Equation 9a:

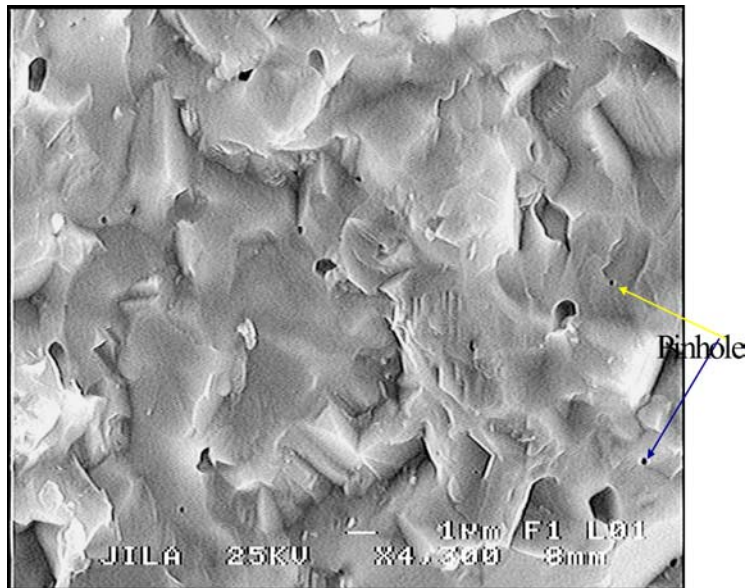
$$\varepsilon_{\text{lattice}} = \frac{d_f - d_0}{d_0} \cdot 100 = \frac{a_f - a_0}{a_0} \cdot 100 \quad (9a)$$

and

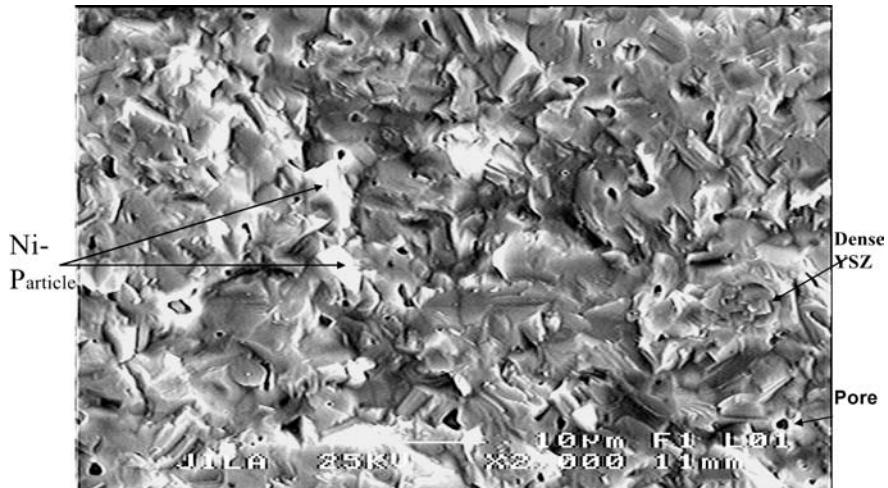
$$\varepsilon_{\text{thermal}} = \bar{\alpha}_{\text{ysz}} \cdot \Delta T \quad (9b)$$

The density of dense YSZ can be accurately measured by Archimedes method to determine the lattice spacing or X-ray diffraction of YSZ can be used to determine lattice parameter and to compute the thermal strains induced in the lattice or vice versa by using Equations 8 and 9.

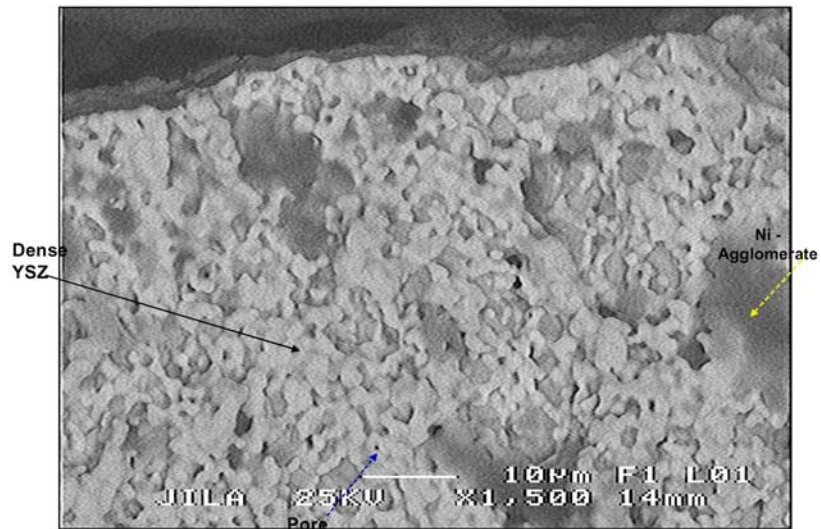
$$\bar{\alpha}_{\text{ysz}} = (8 + 0.003 \cdot T) \cdot 10^{-6}/^\circ\text{C} \quad (10)$$



1-A). Dense 8YSZ



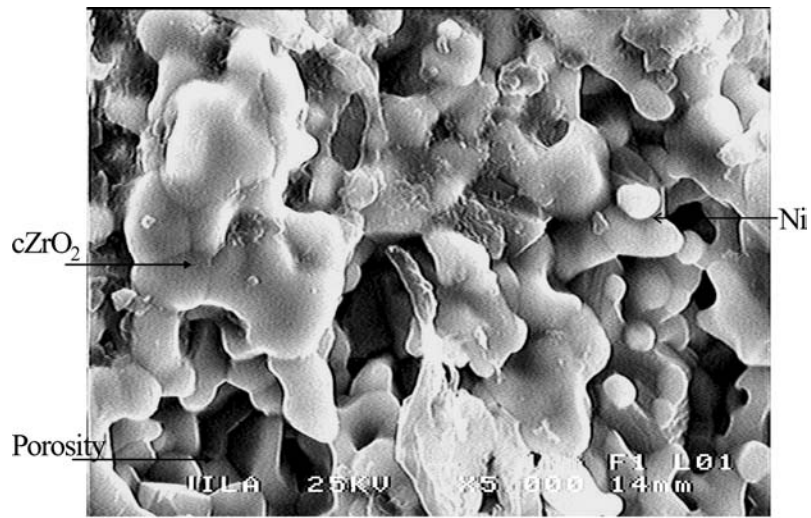
1-B). Porous Ni-YSZ



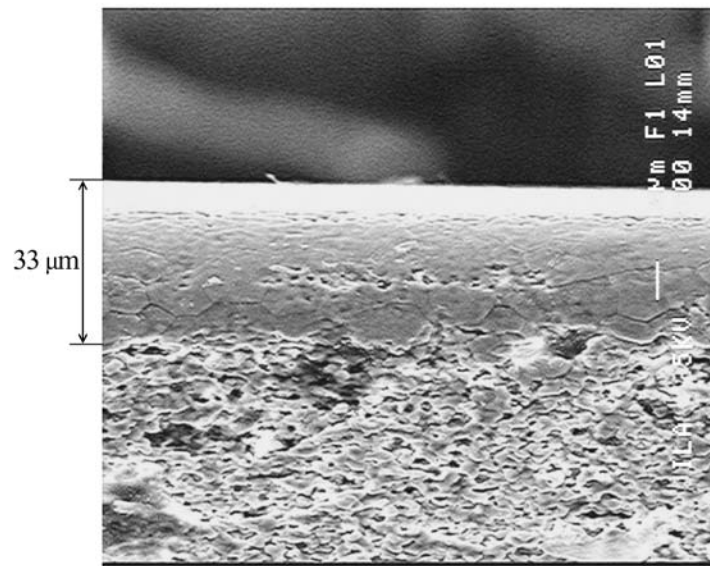
1-C). Back scattered image of phases

Figure 1 (A)–(F). SEM microstructures of dense 8 mole% Yttria Stabilized Zirconia and 25% porosity Ni–ZrO<sub>2</sub>.

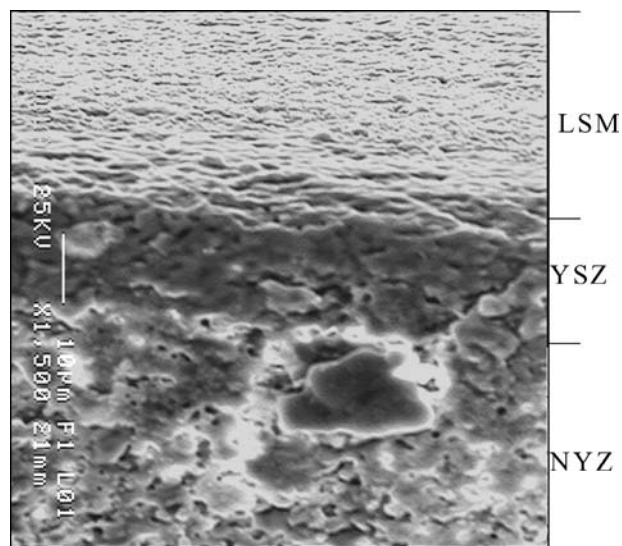
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1-D). Ni-c-ZrO<sub>2</sub> phase at higher resolution.

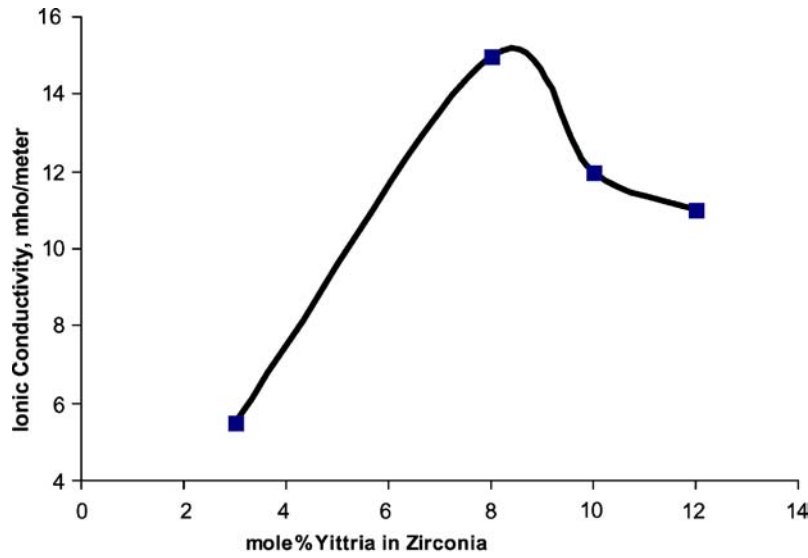


1-E) Bilayer: YSZ film -on-Ni-YSZ substrate.

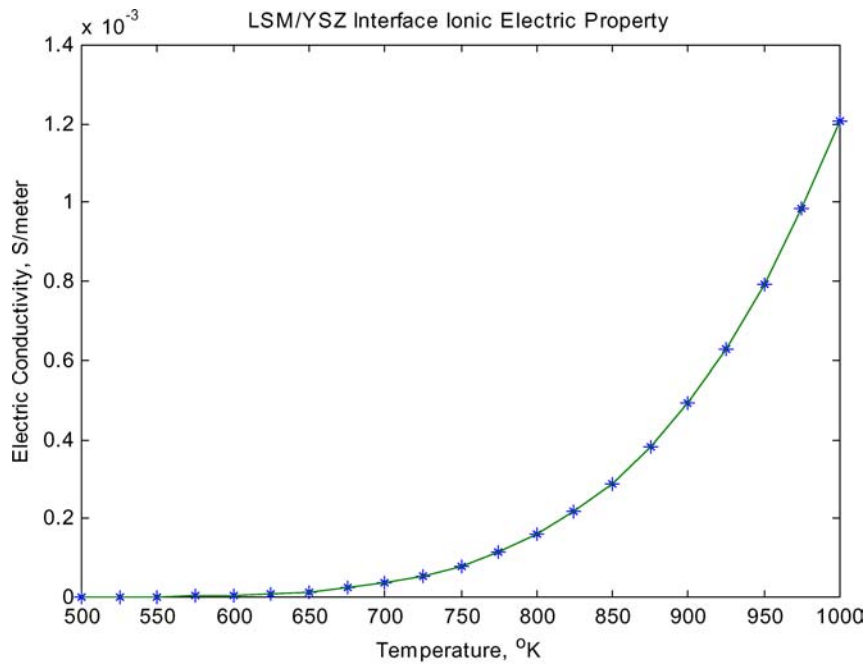


1-F) Trilayer: LSM/YSZ films-on-Ni-YSZ substrate.

Figure 1 (Continued.)



2-A) Ionic conductivity at 1000 °C (data after S.P.S Badwal et. al [27])



2-B). Ionic conductivity of the  $\text{La}_2\text{Zr}_2\text{O}_7$  compounds formed at LSM/YSZ interface.

Figure 2 The ionic conductivity of Yttria Stabilized Zirconia binary system and  $\text{La}_2\text{Zr}_2\text{O}_7$  interface compound conductivity.

Average thermal strains accumulated in bulk materials over a range of temperatures can be estimated based on average coefficient of thermal expansion, CTE of YSZ using Equation (9b). Average CTE of the 8.1 mole% YSZ over temperature ranging from RT~1000°C due to Kandil *et al.* [10], given in Equation (10) can be used to calculate the average thermal strains accumulated over a specified temperature range. The temperature dependant CTE represented by an  $n$ th degree polynomial in Equation (5) can be used to estimate the instantaneous thermal strains. The 6th order polynomial coefficients (Table I) for temperature dependent CTE (ppm/°K) of YSZ due to Sulec *et al.* [3], can be used to calculate the accumulated thermal strains at a given temperature. Residual thermal stress accumulated

in the electrolyte made of YSZ due to thermal mismatch between CTEs of interconnect ( $\alpha_I$ ) and electrolyte ( $\alpha_e$ ) can be calculated by using a polynomial solution [18] for an integral given in Equation 12 based on average constant elastic modulus of electrolyte over a range of temperature extending from stress free ambient ( $T_o$ ) to operating fuel cell temperature ( $T_f$ ).

$$\alpha_{\text{CTE}} f(T) = \sum_{i=0}^n a_i \cdot T^i \quad (11)$$

$$\sigma_{\text{Residual}} = \int_{T_0}^{T_f} E_e \cdot (\alpha_I - \alpha_e) \cdot \Delta T \quad (12)$$

TABLE I Calculated coefficients of nth order polynomial for temperature dependent CTE (PPM/°C) of various fuel-cell materials

Coefficients	$a_0$	$a_1$	$a_2$	$a_3$	$a_4$	$a_5$	$a_6$
Materials							
8YSZ [3]	7.31E-06	-5.58E-09	6.54E-11	-1.57E-13	1.72E-16	-8.94E-20	1.79E-23
Ni-YSZ	8.223	0.0517	-2.00E-04	3.00E-07	-2.00E-10	6.00E-14	
Ni-ZrO <sub>2</sub>	5.4113	0.0833	-0.3E-03	5.00E-07	-4.00E-10	1.00E-13	
E-Brite	9.489	5.40E-03	6.00E-06	4.00E-09			
Interconnect							

Residual thermal stress accumulating from RT ( $T_0=20^\circ\text{C}$ ) to operating temperature of a fuel-cells ( $T_f=800^\circ\text{C}$ ) in YSZ electrolyte for constant high temperature elastic modulus ( $E_e=155$  GPa at  $900^\circ\text{C}$ ) [3] but CTE varying with temperature up to  $800^\circ\text{C}$ , were used in the calculation. The super ferritic E-Brite interconnect with cubic polynomial function of CTE (Equation (11)) varying with temperature is used to predict thermal mismatch between electrolyte and interconnect materials. The residual stress in the electrolyte estimated from spread sheet calculations using Equation (13) was  $\sim 220$  MPa which was lesser than average flexural strength of 277 MPa for Ni-YSZ material determined by three-point bending test. This served as preliminary design stress in the fabrication of five layered unit stack of fuel cell assembly consisting of 60 unit stacks periodically repeating to achieve a target capability  $\sim 1$  KW of electricity.

#### 4.1.2. Elastic properties of dense YSZ

The stiffness constants measured at Iowa State University by Kandil, Greiner and Smith [10] for 11–18 mole% Yttria doped Zirconia are used here to evaluate temperature dependent elastic bulk properties of Yttria Stabilized Zirconia from RT– $800^\circ\text{C}$ . Marked discrepancies in measured stiffness constants are noted from literature for 8 mole% Yttria stabilized Zirconia [8–10, 19–21]. At higher temperatures ( $300$ – $500^\circ\text{C}$ ) measurement of ultrasonic wave propagation in longitudinal or transverse modes during thermal cycling of 8 mole% YSZ was highly unstable and inconsistent [10]. This could be attributed to the minimum dosage of Yttria (8 mole%  $\text{Y}_2\text{O}_3$ ) required for cubic crystal stabilization of Zirconia seems to be unstable for achieving thermodynamic equilibrium at room temperature [10]. Anomalies can be confirmed from oscillatory trend of flexural strength of 8YSZ, reported by More *et al.* [22], which dropped significantly at  $500^\circ\text{C}$  from that at RT and then increased abruptly at  $1000^\circ\text{C}$  may be recalled here for the same thermodynamic instability. The isotropic bulk properties of elasticity of dense Yttria Stabilized Zirconia calculated from section 4.0 using the stiffness coefficients [10] are shown in Figs 3–5. Anisotropy factor of Yttria Stabilized Zirconia defined in Equation (6) clearly increases with molar concentration of Yttria but decreases with temperature for a given dosage of Yttria stabilizer over a range of  $20$ – $700^\circ\text{C}$ . In other words, the

isotropy of Zirconia increases with cubic stabilization by increasing dosage of Yttria (11–18 mole%). Effect of composition on the anisotropy factor of YSZ at  $800^\circ\text{C}$  can be defined by  $AF_{800}=0.0119$  mole%  $\text{Y}_2\text{O}_3 + 0.206$ . Anisotropy factor of 8 mole% YSZ at  $800^\circ\text{C}$  is 0.3012 by extrapolation. At room temperature the anisotropy factor varied from 0.397–0.494 and extrapolated cubic spline fit values of anisotropy at  $800^\circ\text{C}$  varied from 0.341–0.413 with increasing composition of Yttria (11–18%). The Poisson's ratios of YSZ are determined by the ratio of compliance of  $-S_{12}$  to  $S_{11}$  [12]. Poisson's ratio of YSZ increases with composition of Yttria but not significantly affected by temperature up to  $400^\circ\text{C}$ . From  $400$ – $500^\circ\text{C}$ , ratios seem to decrease and converge to fairly constant value thereafter at higher composition. A diverging trend for lower composition (at  $T>500^\circ\text{C}$ ) can't be explained (Fig. 3B). Thermodynamic instability of isotropic cubic YSZ structure for lower Yttria compositions of up to 12 mole% can be realized from lesser significance of Yttria dosage on Anisotropy factor (Fig. 3A). The effect is even more pronounced on the isotropic properties viz, Young's and Shear moduli in Fig. 4 and 5. The isotropic Young's modulus being independent of directional vector ( $R=0$ ) can be determined from Equation 7 by inverse of  $S_{11}$ . The compositional dependence of Young's Modulus of YSZ containing up to 12 mole% Yttria is less significant and have  $E$  values  $\geq 360$  GPa at RT. Young's modulus decreases drastically with temperature by 25–30 GPa from RT to  $700^\circ\text{C}$  for YSZ containing Yttria up to 12 mole%. The isotropic Young's Moduli of YSZ have a maximum value at 12 mole% Yttria compositions over the entire temperature range from  $20$ – $800^\circ\text{C}$ . The isotropic Young's modulus of YSZ decreases thereafter significantly by 25–30 GPa with Yttria composition ranging from 12–18 mole%. Similar trends are observed in Fig. 5, on the compositional and temperature dependence of isotropic Shear modulus of YSZ. The isotropic shear moduli are calculated from half of the inverse of difference in compliance between  $S_{11}$  and  $S_{12}$  [12]. Shear moduli are not significantly affected up to 12 mole% Yttria having  $G \geq 150$  GPa at RT which drops by 10–15 GPa from RT to  $800^\circ\text{C}$ . The compositional dependence of shear modulus of YSZ also peaks at a maximum value 12.1%. The isotropic shear modulus decreases gradually with the composition (12–18 mol% Yttria) and linearly with temperature ( $20$ – $800^\circ\text{C}$ ) by approximately 10 GPa. The bulk moduli of isotropic single crystal of YSZ can

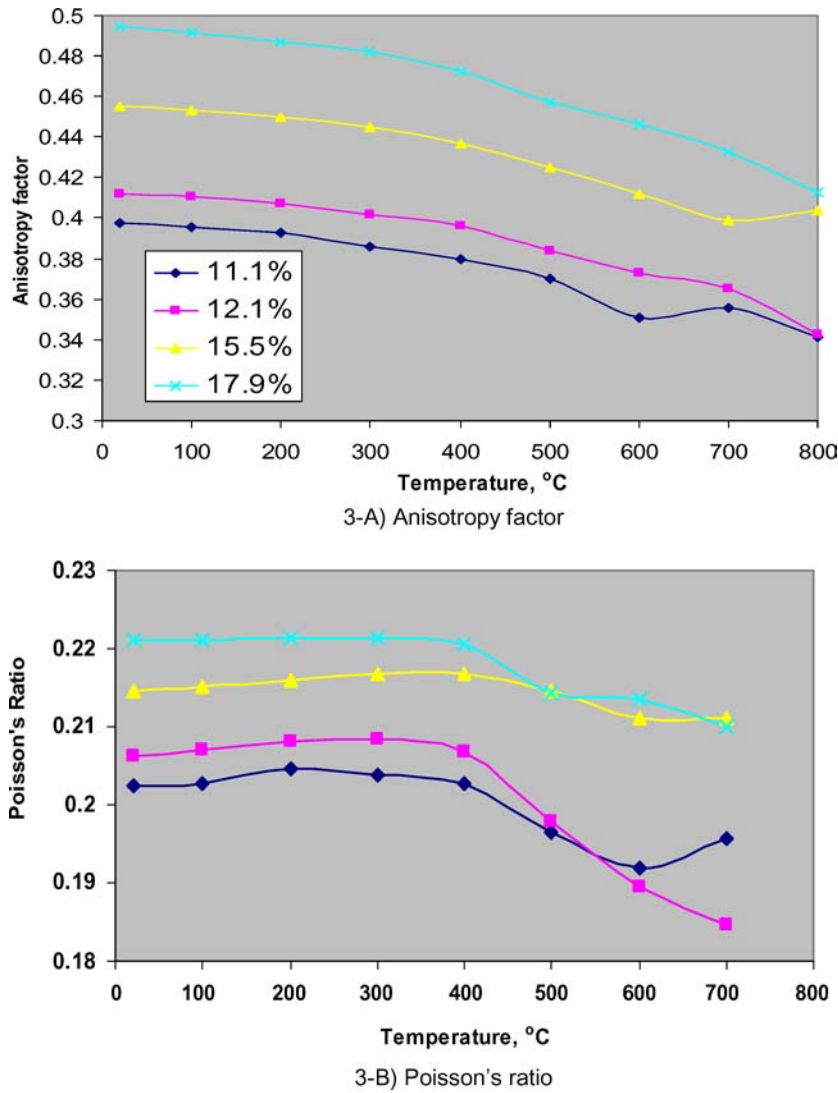


Figure 3 Anisotropy factor and Poisson's ratio varying with temperature and mole% of Yttria doped in Zirconia cubic single crystals

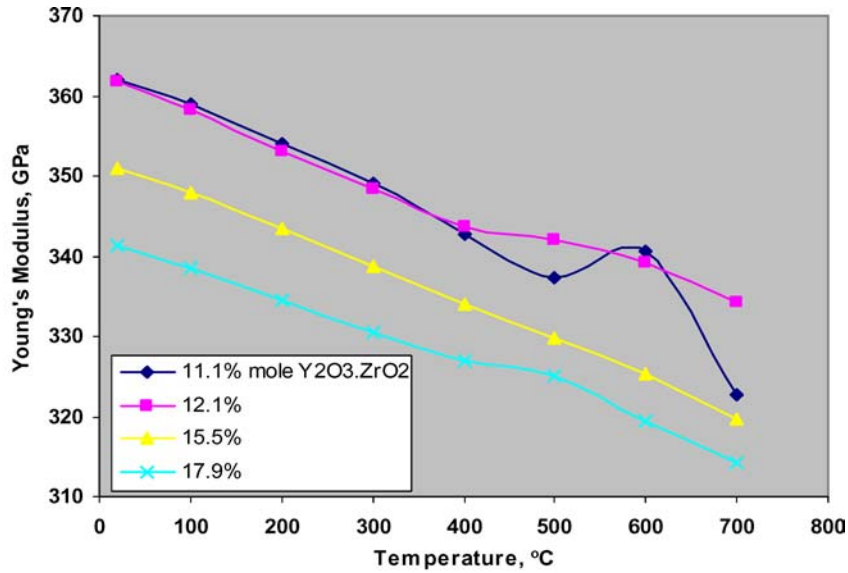
be calculated from one third of the inverse of sum compliance of  $S_{11}$  and  $2S_{12}$ . The cubic polycrystalline elastic moduli are identical to single crystal bulk moduli [10]. With the hypothesis that elastic moduli of polycrystalline cubic YSZ can be derived in terms of stiffness constants by the following relations (13a–c). The polycrystalline properties of various cubic crystal compositions of Yttria Stabilized Zirconia (11–18 mol% Yttria) at room temperatures (20°C) are tabulated (Table II). The polycrystalline Voigt bound average of Young's modulus of 88% c-ZrO<sub>2</sub> ( $E_{100\%ZrO_2} = 220$  GPa) and that of 12% Y<sub>2</sub>O<sub>3</sub> ( $E_{100\%Y_2O_3} = 175$  GPa) has an average Young's modulus of 215 GPa for 12 mole% YSZ which compares with predicted value of 205 GPa for the same within 5% error limit. The predicted elastic modulus for  $E_{\langle 100 \rangle}$  of 12 mole% YSZ is 362 GPa in comparison with actual experimental  $E_{\langle 100 \rangle}$  value of 370 GPa reported for 12 mole% YSZ [23]. The polycrystalline Young's moduli and Shear moduli indicate maximum elastic properties for 12 mole% YSZ until 400°C. By and large, polycrystalline proper-

ties decrease with composition above 12 mole% Yttria and temperatures (RT–800°C). However at temperatures above 400°C, polycrystalline properties of 15.5 mole% Yttria seem to dominate over those of 12 mole% Yttria Stabilized Zirconia.

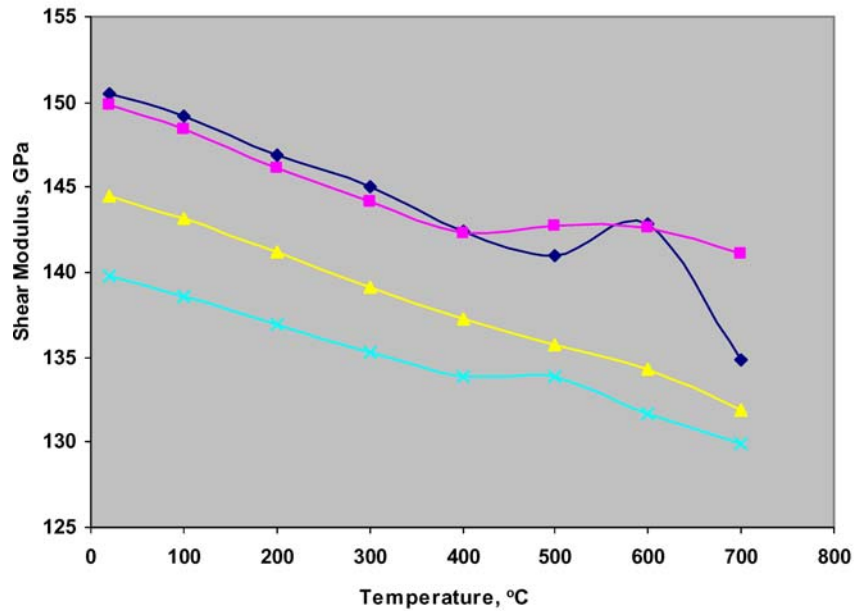
TABLE II Polycrystalline properties of YSZ: Elastic moduli (GPa) and Poisson's ratio

mol% Y <sub>2</sub> O <sub>3</sub>	$E_{YSZpoly}$	$G_{YSZpoly}$	$K_{YSZpoly}$	$\nu$
(a) $T=20^\circ\text{C}$				
11.1	202.77	84.32	113.56	0.202412
12.1	205.23	85.07	116.47	0.206309
15.5	204.93	84.37	119.65	0.214540
17.9	204.00	83.53	121.89	0.221069
(b) $T=700^\circ\text{C}$				
11.1	176.73	73.90	96.82	0.19576
12.1	176.66	74.57	93.36	0.18462
15.5	184.27	76.08	106.25	0.21095
17.9	180.60	74.63	103.76	0.20992





4-A) Young's moduli vs mole% Yittria and Temperature, °C



4-B) Shear moduli vs mole% Yittria and Temperature, °C

Figure 4 Temperature dependence of Isotropic Young's modulus and Shear Moduli of Yittria doped Zirconia.

$$E_{\text{polycrystal}} \equiv K_{\text{singlecrystal}} = \frac{(C_{11} + 2C_{12})}{3} \quad (13a)$$

$$G_{\text{polycrystal}} = \frac{(C_{11} + C_{12})}{6} \quad (13b)$$

and

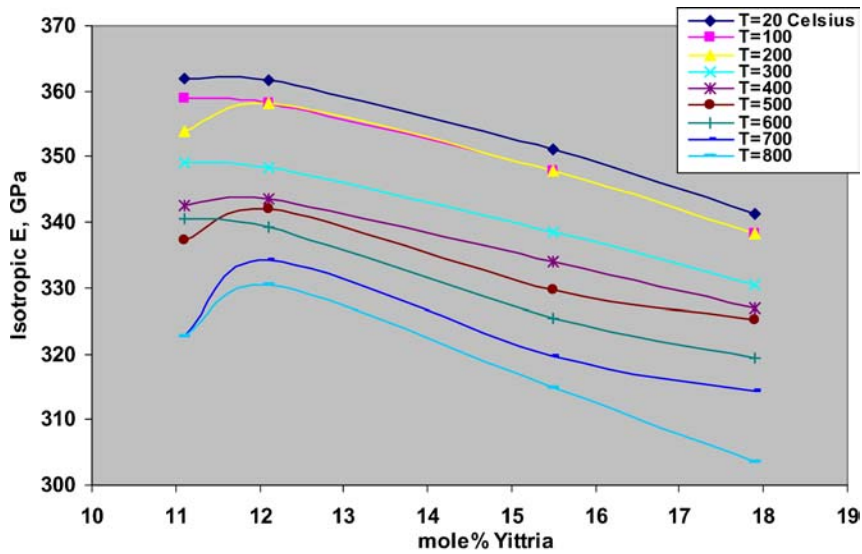
$$K_{\text{polycrystal}} = \frac{(C_{11} + 2C_{12})(C_{11} + C_{12})}{9(C_{11} + C_{12})} \quad (13c)$$

On the basis of the isotropic Young's and Shear moduli as well as polycrystalline properties, it can be inferred that the 12 mole% Yittria Stabilized Zirconia is a better substitute for 8 mole% YSZ from the view point of

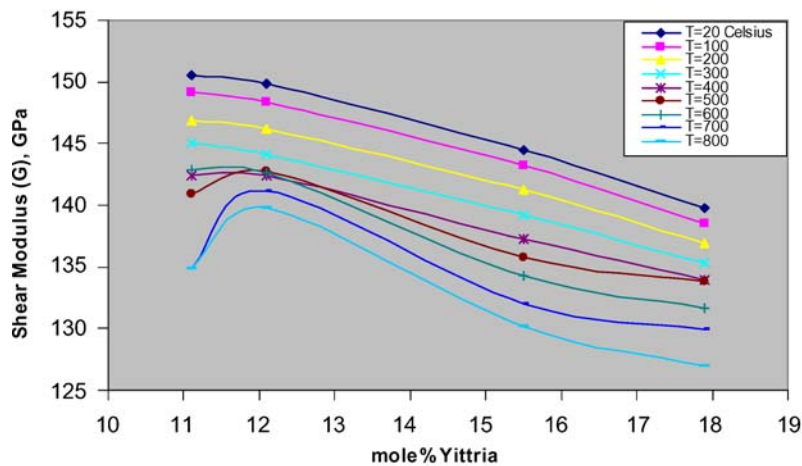
structural, thermodynamic stability, improved isotropic elastic, thermo-elastic mechanical properties and strength required in fuel cell applications.

#### 4.1.3. Porous Ni/YSZ

The microstructure of NiO-YSZ of sintered and reduced to Ni-YSZ is shown in Figure 1B-D, has average porosity of about 25%, being widely dispersed in the matrix. The NiO is reduced to bright shiny Ni particles or agglomerates are found scattered in the matrix of c-ZrO<sub>2</sub>. The spherical Ni particles were seen in microstructure at higher magnification. At times the Ni was found surrounding the rim of pores in matrix. The SEM back scattered images con-



A) Isotropic Young's moduli



B) Isotropic Shear moduli.

Figure 5 Determination of optimal composition of Yttria Stabilized Zirconia for improved thermo-elastic properties based on: A) Young's moduli and B) Shear moduli.

firmed the formation of Ni particles/agglomerates (grey) and cubic  $Y_2O_3 \cdot ZrO_2$  dense (white) structures in reduced atmosphere. Average compositions of these structures were approximately YSZ  $\sim 48\%$ , Ni  $\sim 27\%$  and porosity  $\sim 25\%$ . The porosity manifested into a wide variety of crack morphology in the reduced Ni-YSZ microstructure. Typical types of crack morphology with their average size were circular ( $1-2 \mu m$ ), elliptical ( $1 \times 3 \mu m$ ), pinhole ( $0.5-1 \mu m$ ), or slender slit ( $3-5 \mu m$ ), type of porosity were delineated in reduced Ni-YSZ. However there were no channel or mud cracks [3, 6] seen at the interface of Bilayer (YSZ thin film on Ni-ZrO<sub>2</sub> substrate) or Trilayer (LSM/YSZ film on Ni-ZrO<sub>2</sub> substrate) structures (Fig. 1-E and 1-F). The multilayer films were intact and rigidly held on Ni-ZrO<sub>2</sub> substrate and were free from bending or warpage.

Average thermal expansion (CTE) of Ni/YSZ can be best described by the 5th order polynomial coefficients in Table 1 for Equation 12. A thermal mismatch due to

CTE of the Dense YSZ and porous Ni-YSZ indicates that between the two ( $2-3 \text{ ppm/C}$ ) might result in thermal stress concentration at interface of anode/electrolyte. The excessive stress concentrations could cause delamination MEAs at the interface affecting overall functionality of fuel-cell [6, 7].

#### 4.1.4. Elastic properties of Ni/YSZ

The porosity has a profound effect on the thermo-elastic properties of ceramic cermets. Even though the pores deteriorate the thermo-mechanical properties, porosity improves the electrical performance by ionic transport and thermal diffusivity by local convection of heat. Elastic properties of the porous Ni-YSZ substrates are determined based on elastic properties of the dense YSZ in the previous section. A simple form of Equation (14) can be used to correlate the Young's modulus of the dense and porous

materials of YSZ [24, 26, 27].

$$E_{\text{NYZ}} = E_{\text{YSZ}} (1 - b \cdot P) \quad (14)$$

where,  $P$  — %porosity in NYZ and  $b$  constant  $\sim 3$  for Zirconia based binary system. Young's modulus reduced Ni-YSZ with 25% porosity has  $E_{\text{NYZ}}=54$  GPa using  $E_{\text{YSZ}}=204$  [3, 6] for 8YSZ at the room temperature. Elastic modulus,  $E_{\text{NYZ}}=51$  GPa closely conforms to the reported value of the NYZ  $E_{\text{NYZ}}=54$  at RT [3, 6]. The Poisson's ratio of porous NYZ ( $\mu_{\text{NYZ}}$ ) relating to that of dense YSZ ( $\mu_{\text{YSZ}}$ ) with relatively pores free structure can be calculated using a relation between dense and porous ceramic materials due to Ramakrishnan and Arunachalam [28] in Equation (15): The Poisson's ratio of porous NYZ calculated using (16) and value of  $\mu_{\text{YSZ}}=0.315$  for dense YSZ is found to be  $\mu_{\text{NYZ}}=0.2886$  in comparison with reported value of 0.25 for the same at RT [3]. Thus it is possible to determine isotropic elastic properties porous of Ni-YSZ for various compositions (8–18 mole% YSZ) as a function of temperature (RT–800°C) using Equation 14 and 15.

$$\mu_{\text{NYZ}} = 0.25 \left( \frac{4\mu_{\text{YSZ}} + 3P - 7\mu_{\text{YSZ}} \cdot P}{1 + 2P - 3\mu_{\text{YSZ}} \cdot P} \right) \quad (15)$$

## 5. Conclusion

Structure property correlation of dense and porous Yttria Stabilized Zirconia are established in order to investigate the impact of composition and temperature variation on thermo-elastic properties of isotropic single crystals and the polycrystalline properties. Anisotropy of the cubic structure is defined to understand the compositional variations in the dense YSZ. Elastic bulk properties are numerically predicted using the elastic stiffness constants reported in the literature. Thermo-elastic properties are evaluated for 8 mole% and 12 mole% Yttria Stabilized Zirconia. Effect of Yttria dosage in Zirconia on its lattice parameter, lattice constant, thermal diffusivity and ionic conductivity of the Yttria Stabilized Cubic Zirconia are investigated. A dosage of 8 or 9 mole% of Yttria in Zirconia seems to improve the ionic conductivity of the stabilized c-ZrO<sub>2</sub>. The ionic conductivities of 8 and 12 mole% YSZ are not significantly altered at intermediate temperatures (600–800°K) however the former has slightly better ionic conductivities at higher temperatures. The ionic conductivity of La<sub>2</sub>Zr<sub>2</sub>O<sub>7</sub> type of compound formed in LSM/YSZ type of composite cathodes is found to be 1.2E-03 mho/meter at 1000°K. Even though a minimum of 8 mole% Yttria is necessary to retain the cubic crystal structure at room temperature in Zirconia based material it may be insufficient to achieve the thermal stability and thermo-dynamic equilibrium of Yttria Stabilized Zirconia. A numerical prediction of isotropic elastic properties have indirectly confirmed suspected suscepti-

bility of 8 mole% Yttria Stabilized Zirconia to thermal and harmonic instabilities during thermal cycling.

Poisson's ratio increases with composition but remain constant up to a temperature 400°C. The compositional effect is not well defined at higher temperature. Poisson's ratio decreases with temperature from 400–800°C. Poisson's ratio of porous Ni-YSZ can be determined for any given temperature and composition from Equation (15). The isotropic elastic and shear Moduli increase with the composition of Yttria up to 12 mole% and decrease with any further increase of Yttria composition. Optimal isotropic elastic properties correspond to a composition of 12.1 mole% Yttria in cubic Zirconia. The directional dependence of single crystal elastic modulus for (100) orientation was found to be 362 GPa compared to experimental value of 370 GPa [23]. Elastic properties of porous Ni-YSZ can be determined from Equation 17 and elastic data of dense YSZ.

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